

HIGGS AT e^+e^- LINEAR COLLIDERS

(Theor. Part)

1. INTRODUCTION

- general physics scenario
- 2-phase LC project

2. THE BASE: SM HIGGS MECHANISM

- study Higgs mechanism for EW SB
- precision analysis: looking beyond SM

3. SM EXTENSION: SUPERSYMMETRY

- MSSM as a framework
- generalized scenarios

4. ALTERNATIVE: STRONG EW SYMMETRY BREAKING

- composite Higgs particles
- "no" Higgs

5. CONCLUSIONS

1. INTRODUCTION

LC target: precision analysis of Higgs mechanism / elwSB
 ↓
 telescope for exploring new physics areas

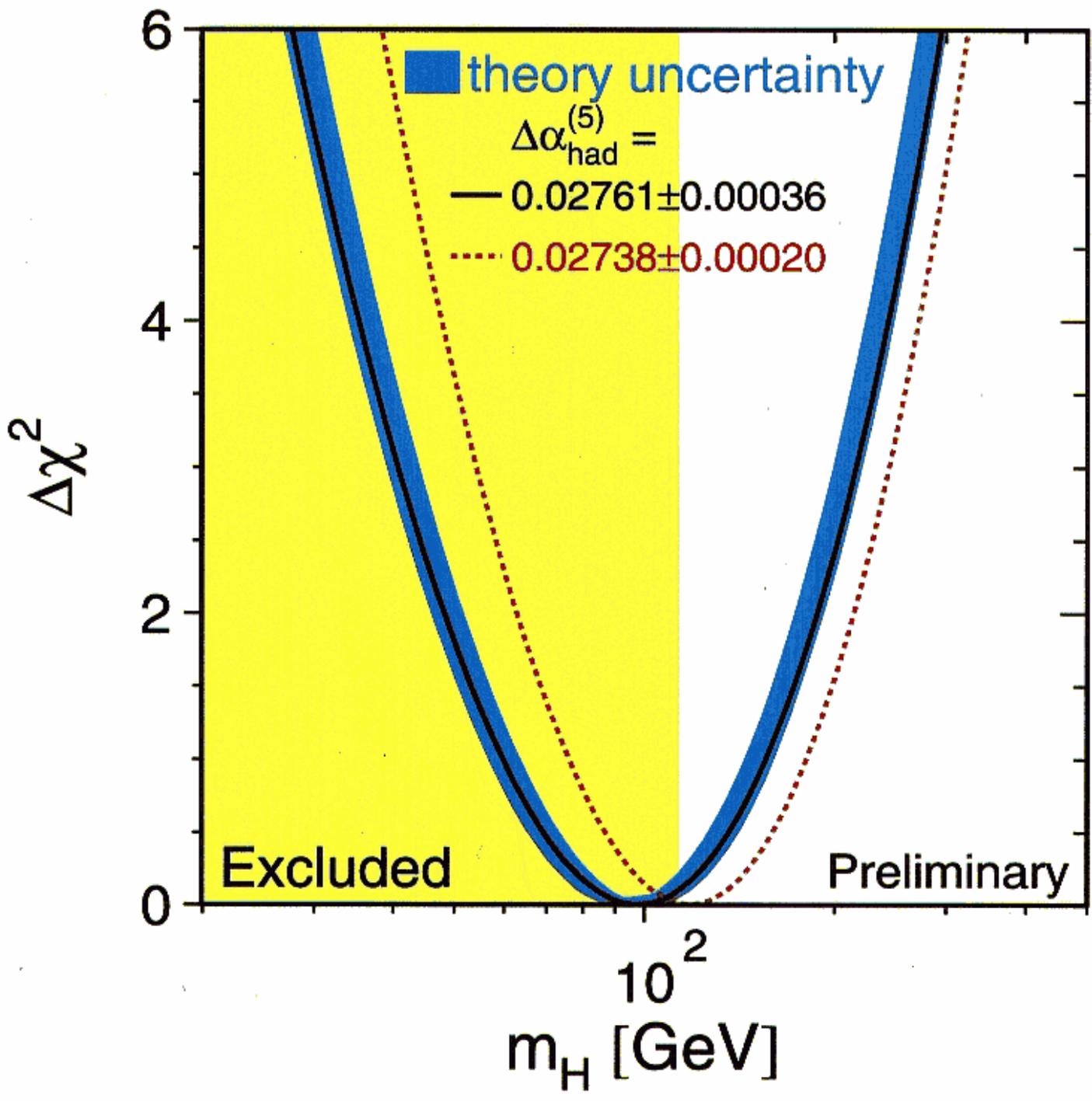
Scales:

- $v \approx 246 \text{ GeV}$ elwSB scale: energy range $\approx \mathcal{O}(1 \text{ TeV})$
 expected crucial for elwSB
- Λ_* scale of new [mg] interactions:
 - $\Lambda_* \approx M_{\text{PL}}$: light fundamental Higgs boson
 naturalness: SM \rightarrow SUSY at $\mathcal{O}(\text{TeV})$
 - $\Lambda_* \ll M_{\text{PL}}$: composite Higgs particle [intm. scale]
 "no" Higgs: strong WW interaction

LC PROJECT:

		\sqrt{s}	$\int \mathcal{L} \text{ run}$	
Phase I	TESLA	$M_Z \dots 500 \dots 800 \text{ GeV}$	$300 \dots 500 \text{ fb}^{-1}$	TDR
	JLC/NLC	$M_Z \dots 500 \dots 888 \text{ GeV}$	$\dots 200 \dots 340 \text{ fb}^{-1}$	
Phase II	CLIC	$\dots 3 \dots 5 \text{ TeV}$	$0.3 / 1 \text{ ab}^{-1}$	2nd g.

ELWVG



theory uncertainty

$$\Delta\alpha_{\text{had}}^{(5)} =$$

$$\text{— } 0.02761 \pm 0.00036$$

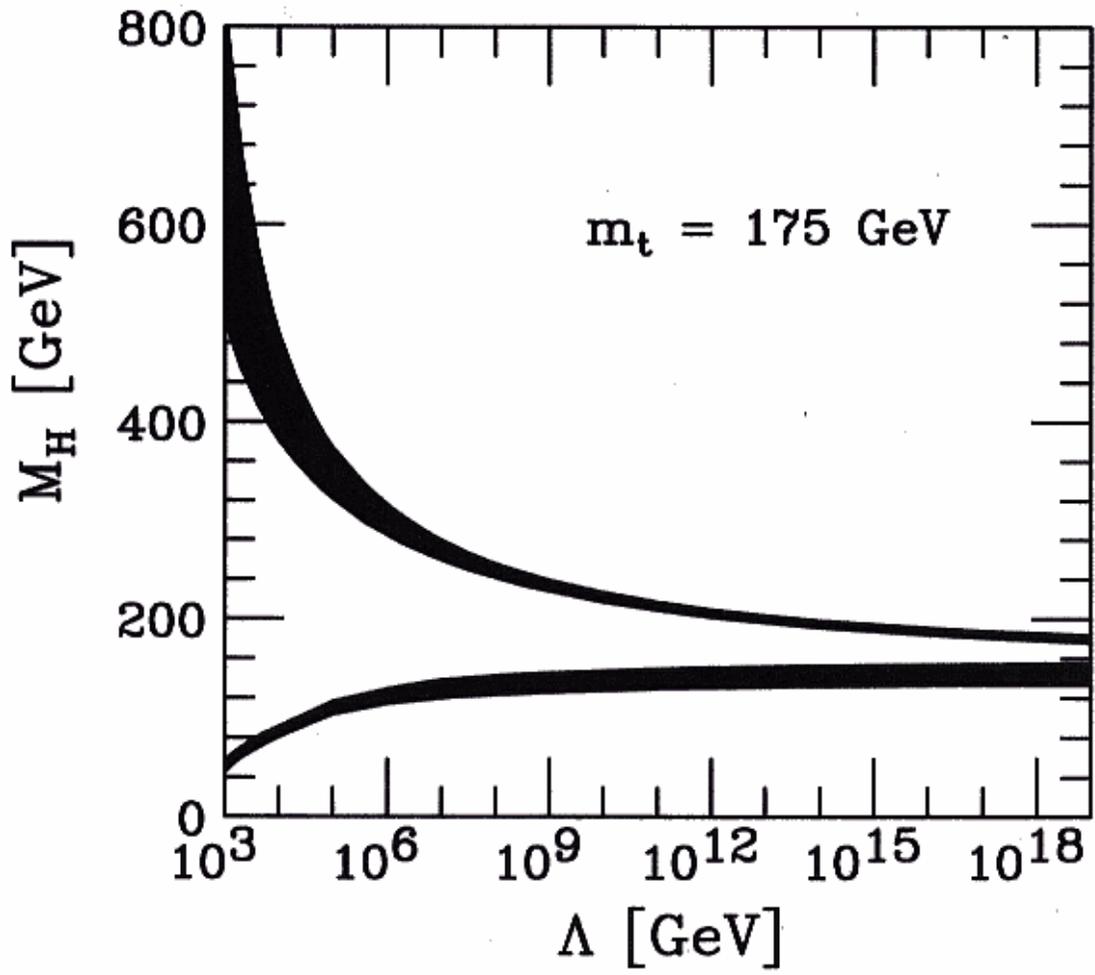
$$\cdots 0.02738 \pm 0.00020$$

Excluded

Preliminary

10^2
 m_H [GeV]

Rienelmann



doubling e^- am : 1.2 TeV

LC $\gamma\gamma$ mode : monochromatic $\Delta \sim 10\%$
luminosity $\sim 1/2 \mathcal{L}_{e^+e^-}$

Giga z : LE operation $\geq 10^9 z$ per a

Project-bandwidth : M_z to several TeV
high luminosity
clean exp. environment

→ high-resolution picture of e^+e^- symmetry breaking

→ extrapolations to scales not accessible directly

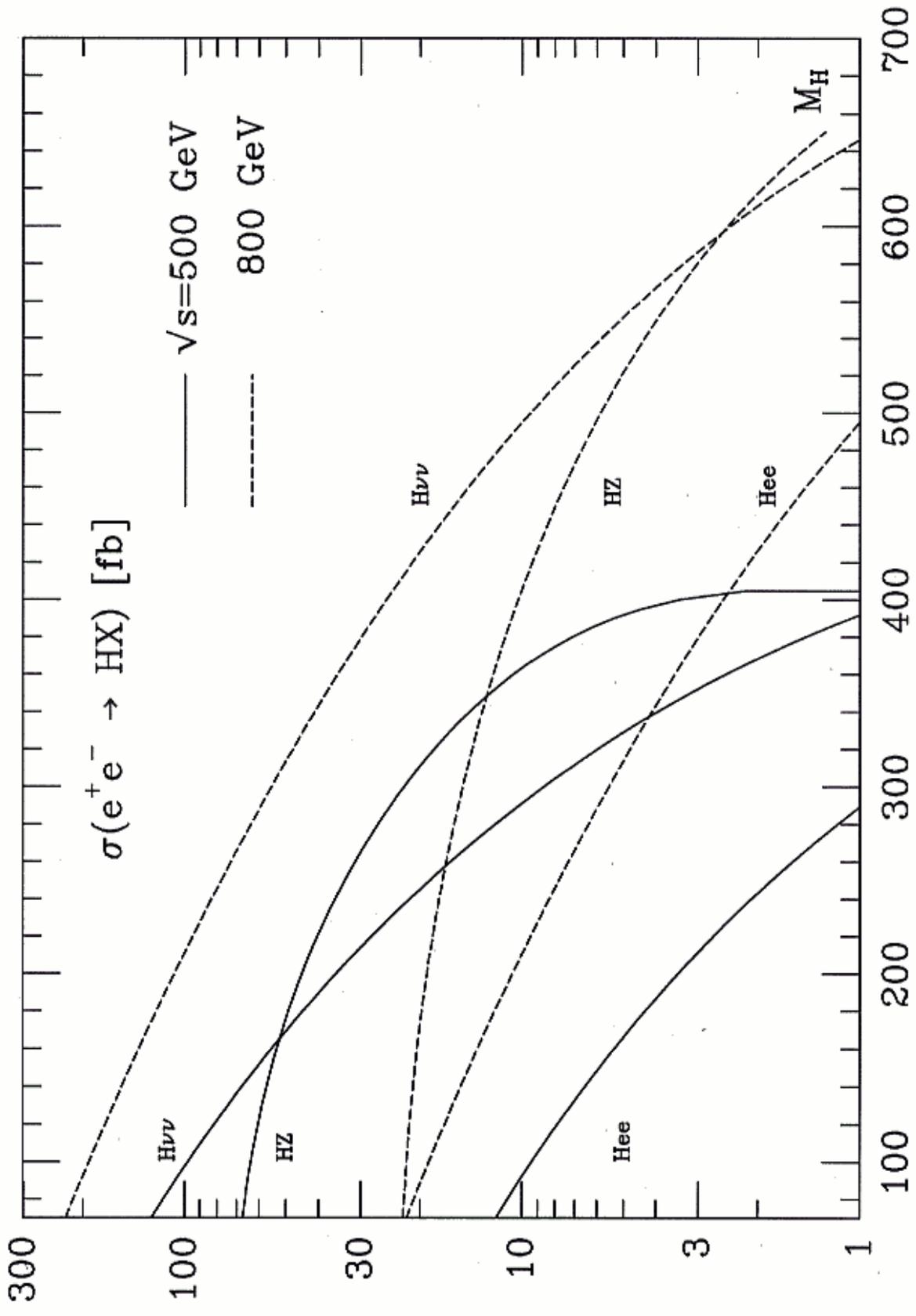
2. THE BASE: SM HIGGS MECHANISM

EW SB in fundamental scalar iso-doublet sector:

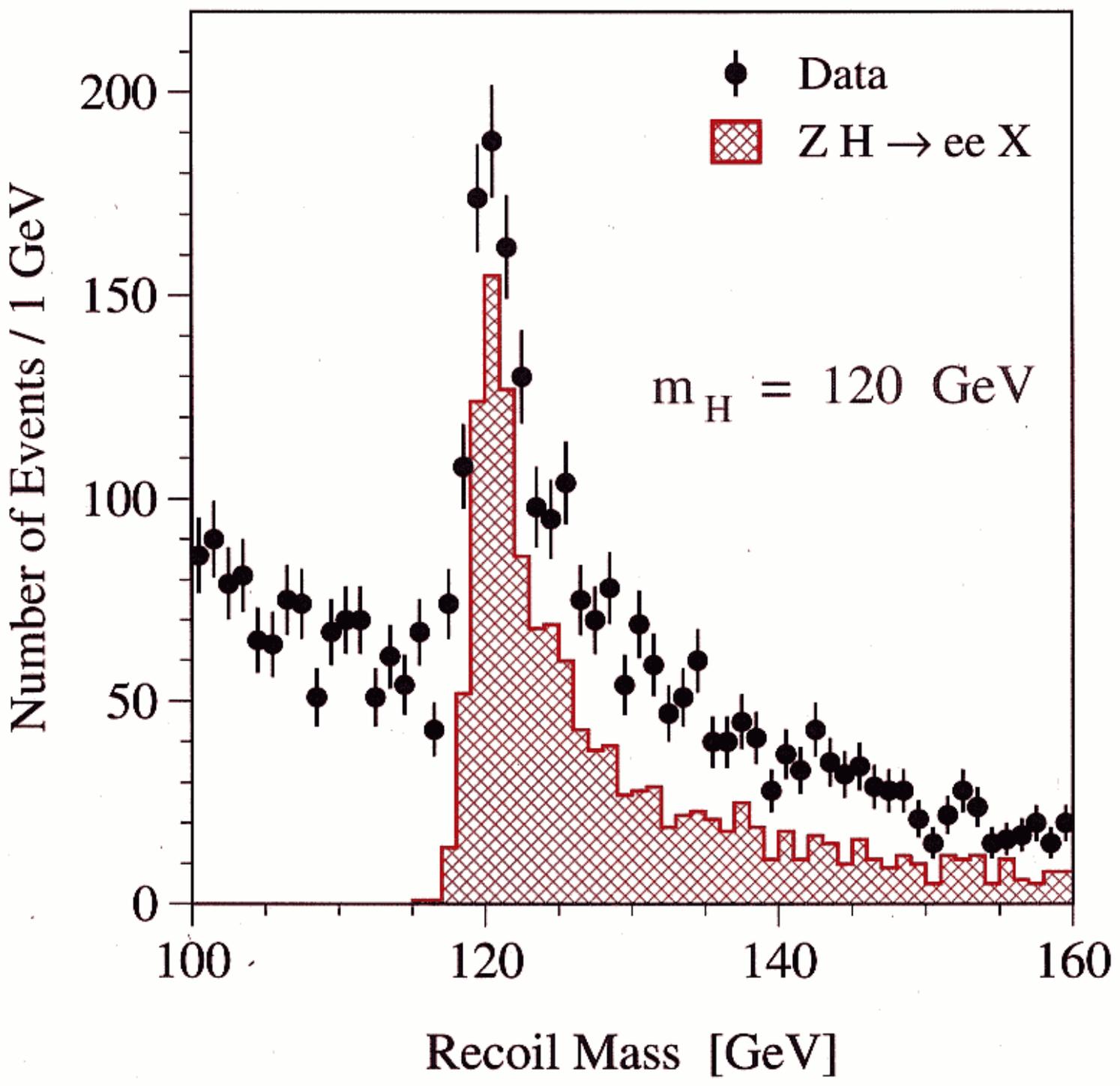
Higgs boson : all properties determined if mass fixed

$M_H \approx 200$ GeV ← analysis of precision exp.

← extdg up to M_{PL}
prerequisite $s_w^2 \sim 0.2$ }



*Garwin - Abri
Lohmann*

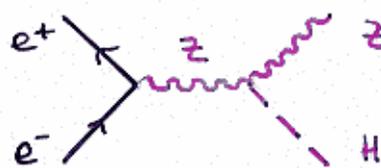


TASK: establish Higgs mechanism experimentally in three steps

- (1) Higgs excitation = Higgs boson must be discovered [LEP2?]
TeVatron
LHC
- (2) generating masses by interaction with Higgs field:
coupling \sim mass [LHC]
LC
- (3) Higgs field in vacuum generated by spont. sym. brk:
reconstruction of Higgs potential LC

LC PRODUCTION OF SM HIGGS BOSON

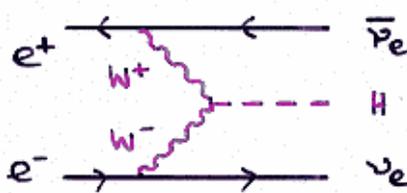
(a) Higgs-strahlung:



$$\sigma = \frac{G_F^2 M_Z^4}{96\pi s} K_5 \beta$$

\sqrt{s} lower part

(b) W fusion:



$$\sigma \approx \frac{G_F^2 M_W^4}{4\sqrt{2}\pi^3} \log \frac{s}{M_H^2}$$

\sqrt{s} upper part

F: $\sqrt{s} \sim 300 \text{ GeV}$ / SR $\sim 500 \text{ fb}^{-1}$

$M_H \sim$ intermediate mass range: 10^5 Higgs bosons

F: almost background-free

Coupling	$M_H = 120 \text{ GeV}$	140 GeV
g_{HWW}	± 0.012	± 0.020
g_{HZZ}	± 0.012	± 0.013
$g_{H\mu\mu}$	± 0.030	± 0.061
g_{Hbb}	± 0.022	± 0.022
g_{Hcc}	± 0.037	± 0.102
$g_{H\tau\tau}$	± 0.033	± 0.048
g_{HWW}/g_{HZZ}	± 0.017	± 0.024
$g_{H\mu\mu}/g_{HWW}$	± 0.029	± 0.052
g_{Hbb}/g_{HWW}	± 0.012	± 0.022
$g_{H\tau\tau}/g_{HWW}$	± 0.033	± 0.041
$g_{H\mu\mu}/g_{Hbb}$	± 0.026	± 0.057
g_{Hcc}/g_{Hbb}	± 0.041	± 0.100
$g_{H\tau\tau}/g_{Hbb}$	± 0.027	± 0.042

Table 2.2.6: Relative accuracy on Higgs couplings and their ratios obtained from a global fit (see text). An integrated luminosity of 500 fb^{-1} at $\sqrt{s} = 500 \text{ GeV}$ is assumed except for the measurement of $g_{H\mu\mu}$, which assumes 1000 fb^{-1} at $\sqrt{s} = 800 \text{ GeV}$ in addition.

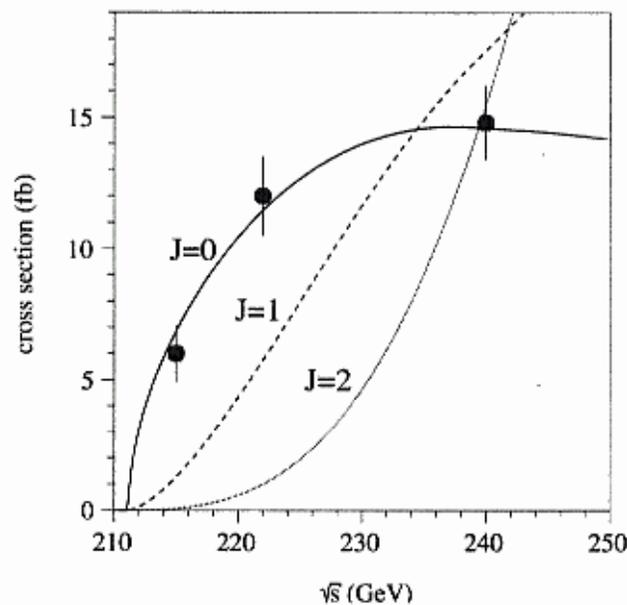
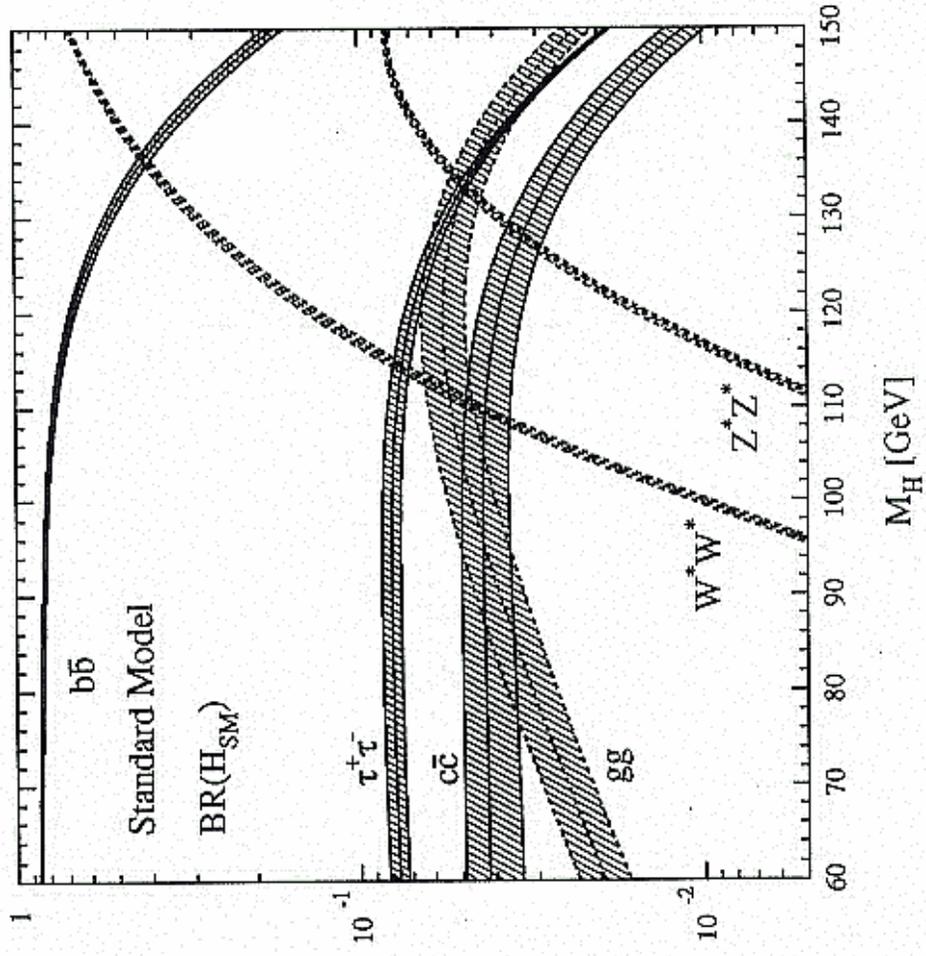
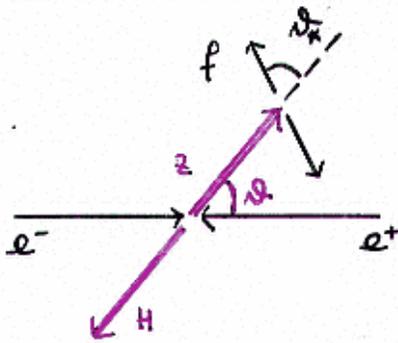


Figure 2.2.7: Simulated measurement of the $e^+e^- \rightarrow H^0Z$ cross section for $M_H = 120 \text{ GeV}$ with $20 \text{ fb}^{-1}/\text{point}$ at three centre-of-mass energies compared to the predictions for a spin-0 (full line) and examples of spin-1 (dashed line) and spin-2 (dotted line) particles.

Djouadi et al.



$J^{PC} = 0^{++}$:



threshold : $\sigma \sim \beta \sim \sqrt{s - (m_H + m_f)^2}$

$\propto [1 + \cos^2 \theta] \cdot [\sin^2 \theta^*]$

v.v.

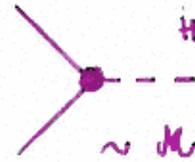
\Rightarrow ruling out : $J^P = 0^-, 1^\pm, 2^\pm, \dots$

CP admix : $H = [0^{++}] + i\eta [0^{+-}]$

$\frac{d\sigma}{d\cos\theta} = \sin^2\theta + 8\eta p \cos\theta \Rightarrow$ sens. $\eta \lesssim 0.03$

HIGGS COUPLINGS

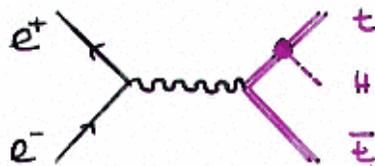
SM masses generated by interaction with Higgs field: $g \sim M$



\leftarrow production : $\sigma(H) \sim g^2 (HZZ)$ etc

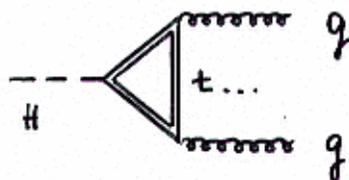
\leftarrow decay : $BR(f) \sim g^2 (Hff)$ etc

\leftarrow top couplg :



Djouadi et al
Dawson et al
Dittmaier et al

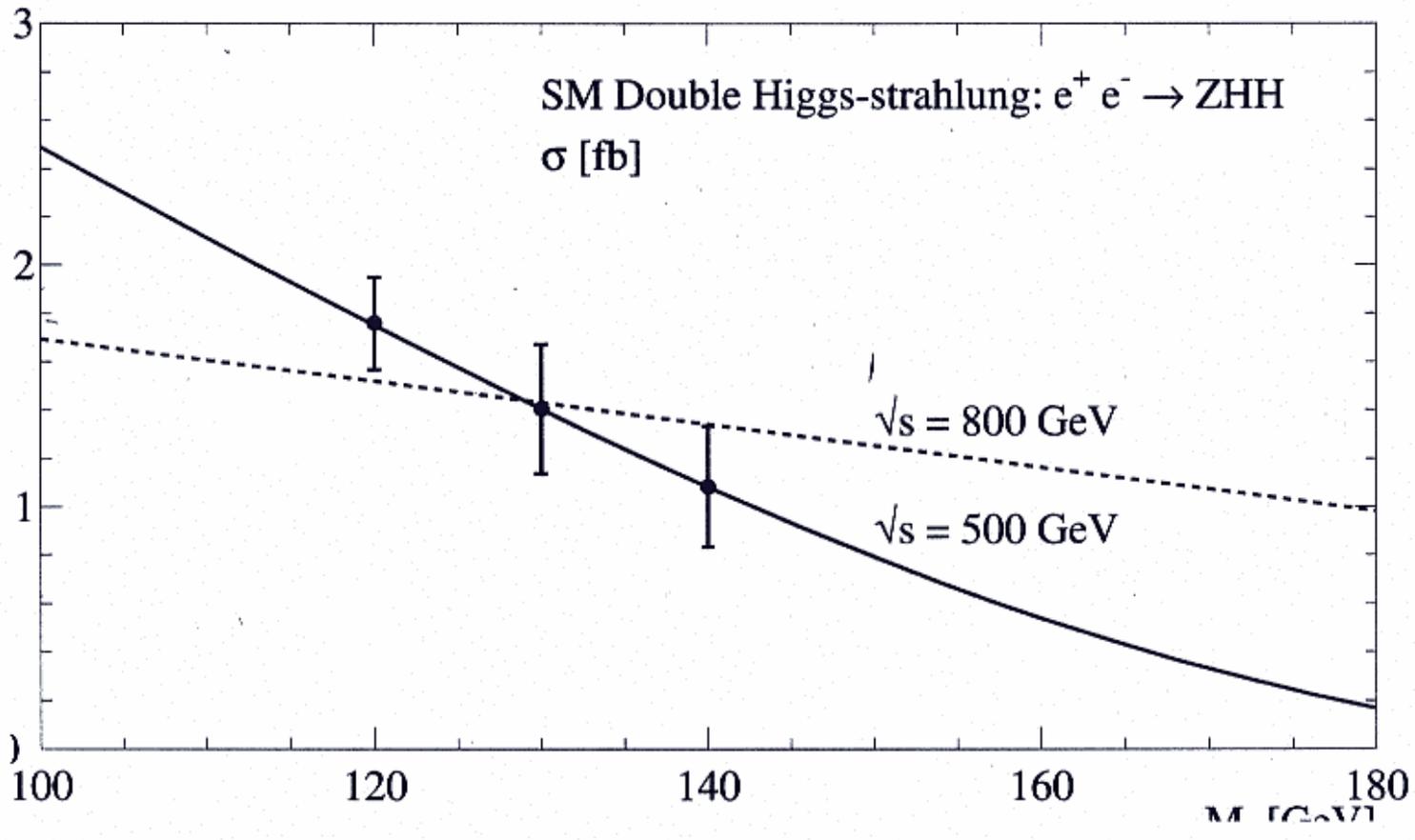
\leftarrow loop cplg :



and $H \rightarrow \gamma\gamma$ decay

$\gamma\gamma \rightarrow H$ production

MÜHLEITNER ER
RAY ER



SUMMARY:

coupling	$M_H = 120 \text{ GeV}$	140 GeV
HWW	0.012	0.020
HZZ	0.012	0.013
Hbb	0.022	0.022
Hcc	0.037	0.102
H $\tau\tau$	0.033	0.048
Htt	0.030	0.061

Bahagia
Dard ea
jeste ea
Baar ea

(i) same mechanism
for gauge bosons
and fermions

(ii) up/down ratio:

$$\frac{c}{b} = \frac{m_c^2(M_H)}{m_b^2(M_H)}$$

(iii) lepton/ quark ratio:

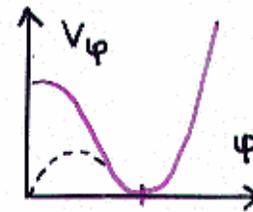
$$\frac{\tau}{b} = \frac{m_\tau^2}{3m_b^2(M_H)}$$

HIGGS POTENTIAL

self-interactions of scalar fields

→ shift of minimum $v/\sqrt{2} \neq 0$

→ electroweak symmetry breaking



$$V = \lambda \left[\phi^2 - \frac{1}{2}v^2 \right]^2 \Rightarrow \lambda_2 = M_H^2/2$$

$$\Rightarrow \lambda_3 = M_H^2/2v$$

$$\lambda_4 = M_H^2/8v^2$$

→ oscill about min.

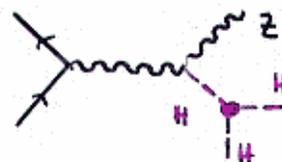
→ damp at $\phi = 0$

→ final form

HHH:

double Higgs-branching:

$$\sigma \sim \alpha^2 g^2 (HHH) \leq 1 \text{ fb}$$



sensitivity

$$\approx 0.18$$

Carreno et al

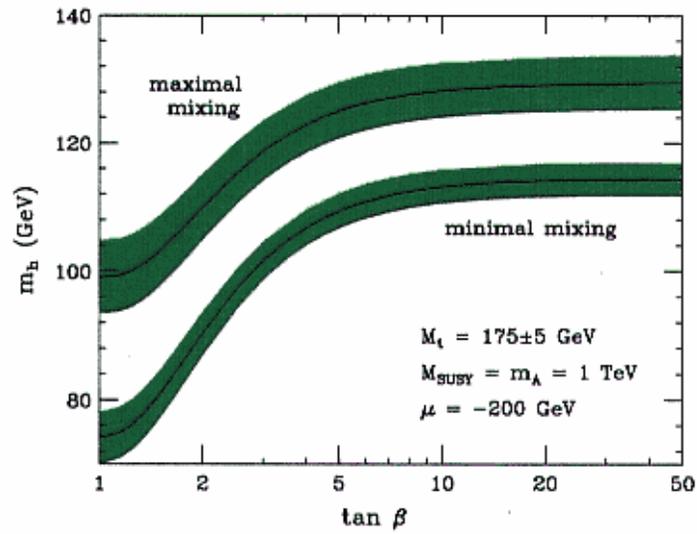
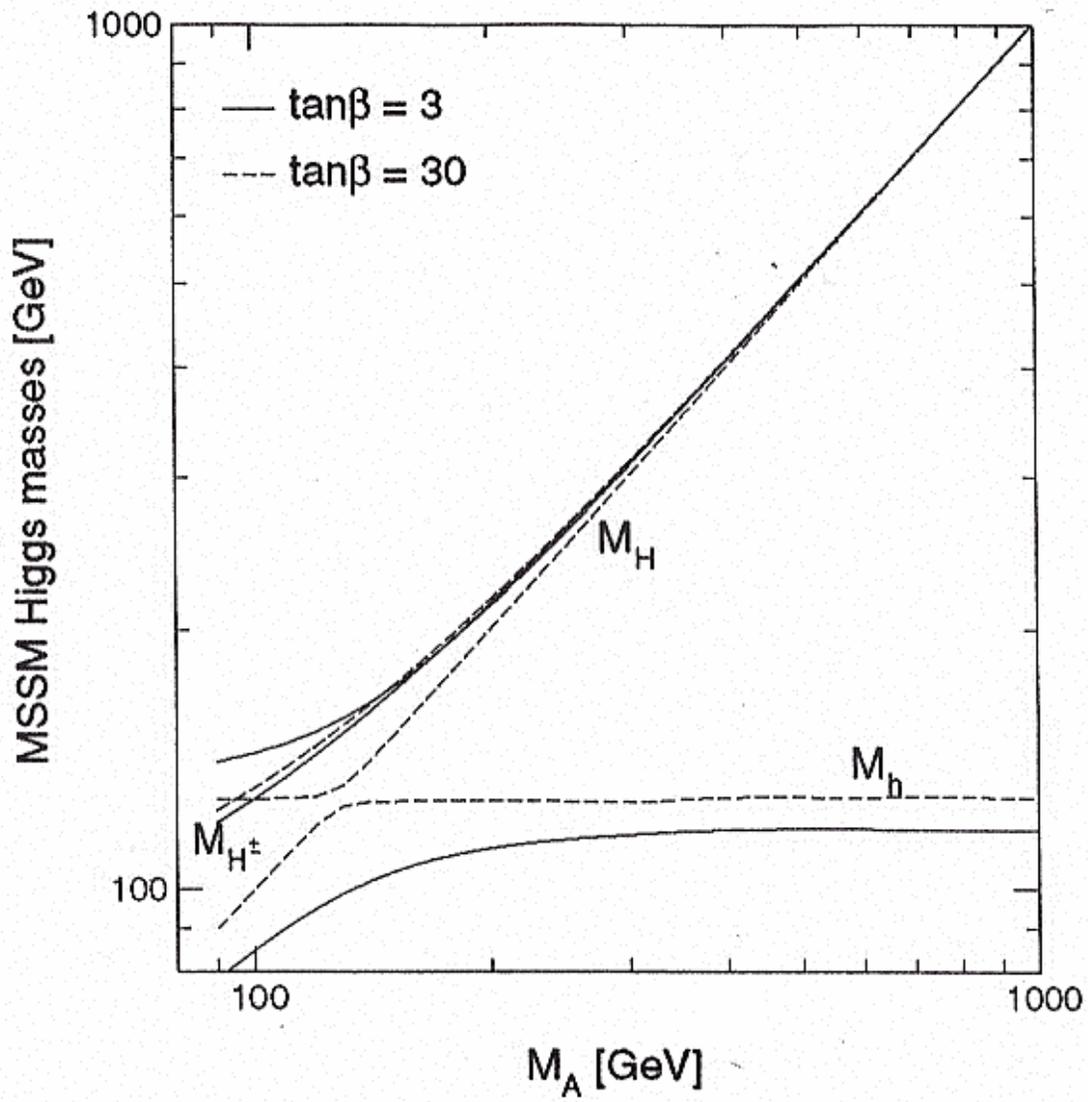


Figure 11: The radiatively corrected light CP-even Higgs mass is plotted as a function of $\tan \beta$, for $M_{\text{SUSY}} \equiv M_Q = M_U = M_D = 1$ TeV and $m_A = 1$ TeV, for the maximal mixing [upper band] and minimal mixing [lower band] benchmark cases [see caption to fig. 10]. The impact of the top quark mass is exhibited by the shaded bands; the central value corresponds to $M_t = 175$ GeV, while the upper [lower] edge of the bands correspond to increasing [decreasing] M_t by 5 GeV.



HIGH PRECISION :

a) establish exp. ehw SB in SM etc: *valuta physica sui generis*

b) consequences BSM examples:

■ new strong interactions at intermediate scale:

changing fermion cplgs, self-interactions, ...?
affecting loop couplings?

Snowman

■ only top mass generated by Higgs mechanism,
other fermion mass "radiatively":

\rightarrow bb

\rightarrow gg dominant for low Higgs masses

Fitzedon

■ Xtra dimensions:

new scalars mixg with Higgs boson

RS: radion \leftrightarrow Higgs

L_{gg} L_{bb} : change BR(H)

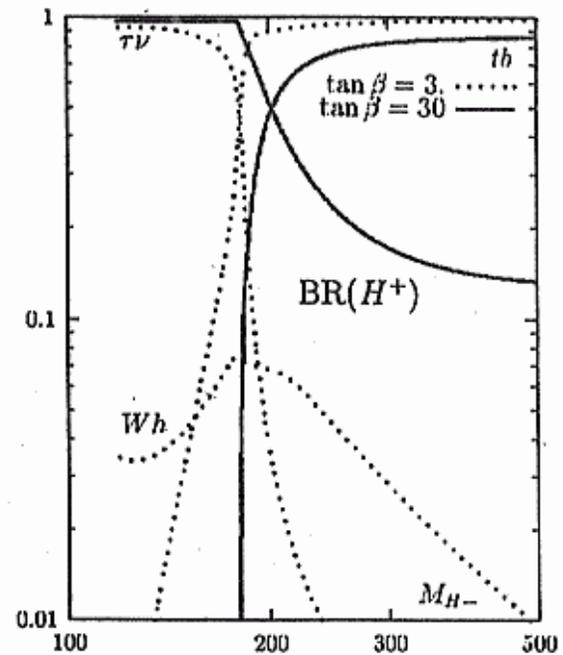
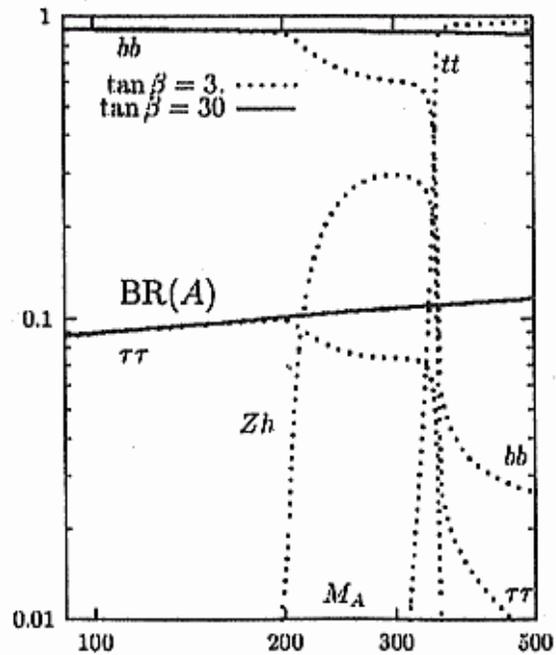
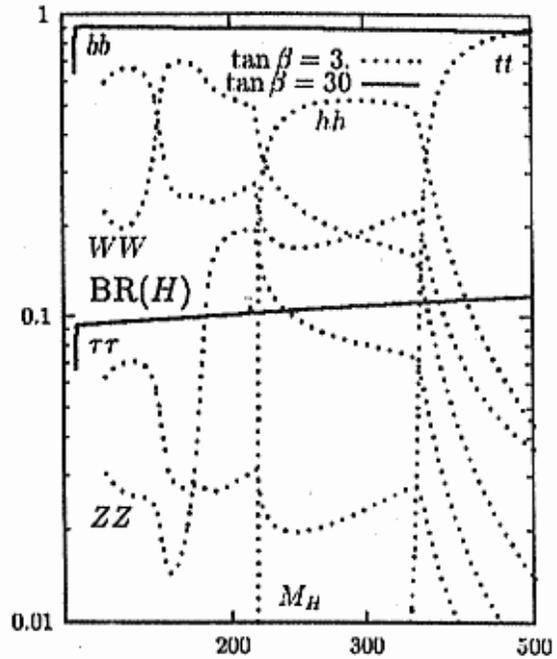
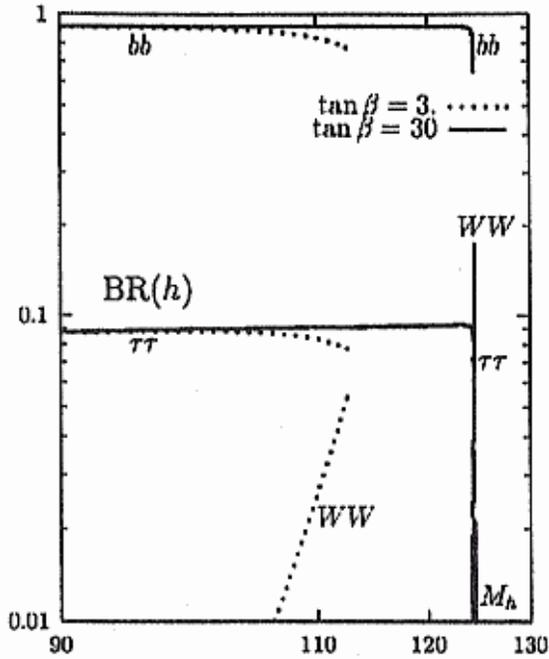
*Csaki ea
Cheung*

3. SUPERSYMMETRY

MSSM: 2 doublets \rightarrow spectrum of 5 states:

h^0 \approx 135 GeV
 H^0, A^0, H^\pm \sim $O(v)$ to $O(\text{TeV})$ } parameters:
 $M_A, \tan\beta$

\ large: degen



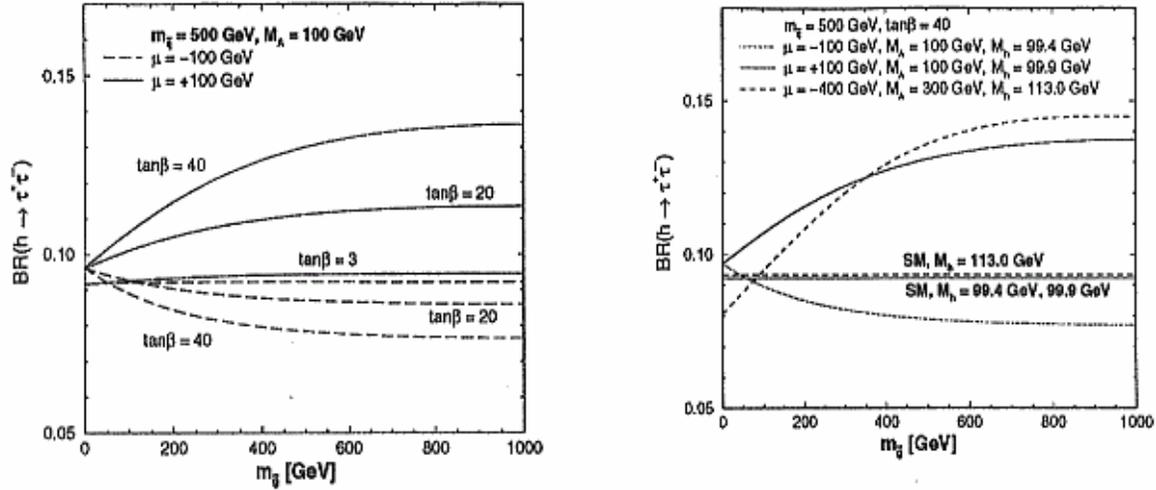
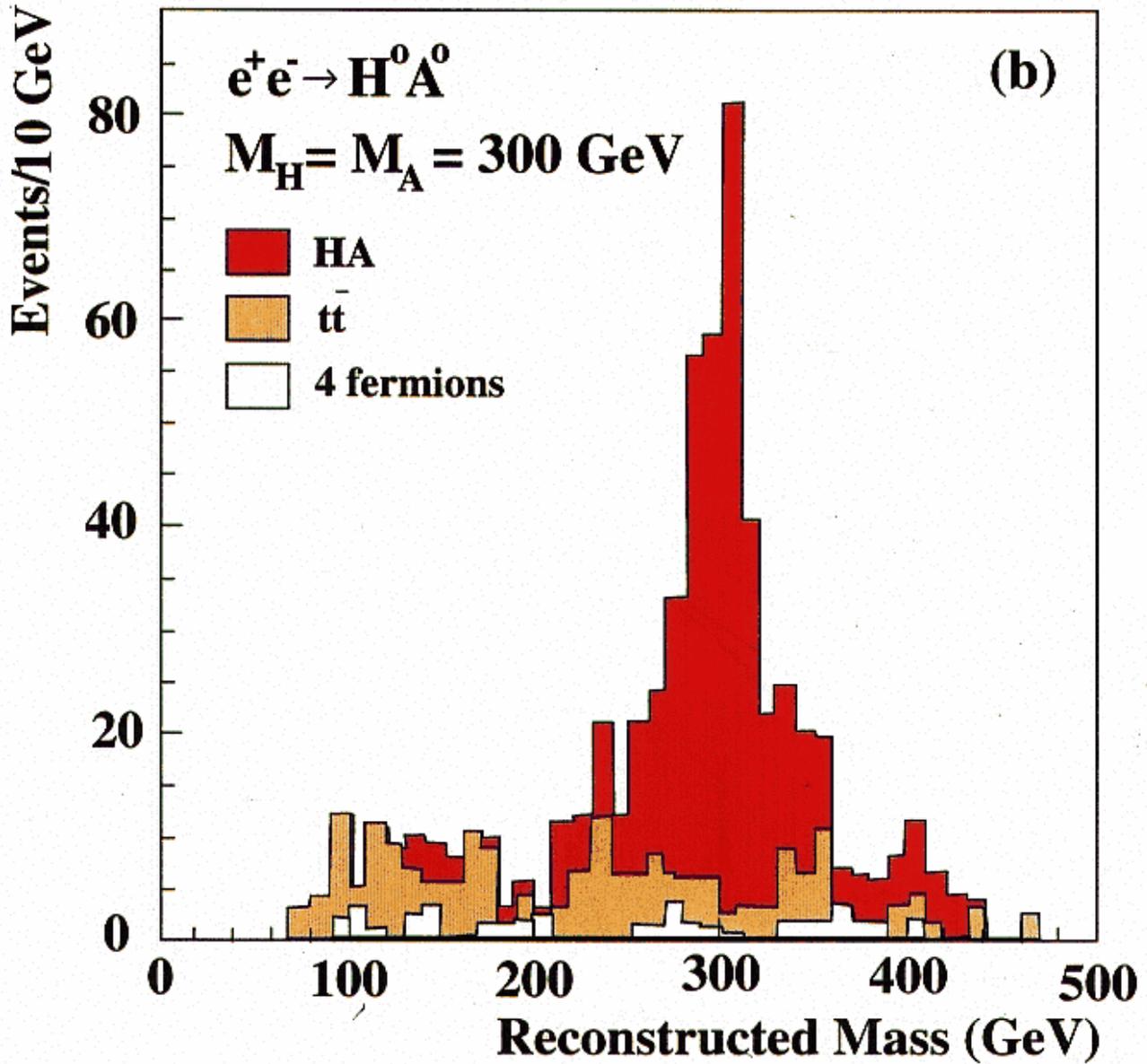


Figure 7: $BR(h \rightarrow \tau^+ \tau^-)$ is shown as a function of $m_{\tilde{g}}$. The left plot corresponds to three different values of $\tan\beta$ and $\mu = \pm 100$ GeV. The right plot corresponds to $\tan\beta = 40$ and the two scenarios $M_A = 100$ GeV, $\mu = \pm 100$ GeV and $M_A = 300$ GeV, $\mu = -400$ GeV. The other parameters in both plots are $m_{\tilde{g}} = 500$ GeV, $X_t = 500$ GeV, $M_A = 100$ GeV, $M_2 = 500$ GeV, $A_b = A_t$. The Higgs boson masses given in the legend are averaged masses over the interval $0 < m_{\tilde{g}} < 1000$ GeV where the variation of M_h is about ± 1 GeV for the above considered parameters. The horizontal lines represent the SM result using the corresponding Higgs boson masses.

Heinzeper a



DECAY MODES:

- Higgs → SM particles
- cascade decays
- supersym. particles

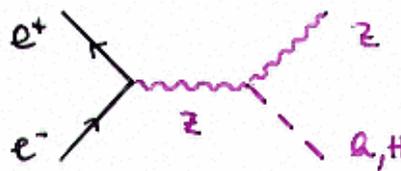
- $h^0 \rightarrow b\bar{b} \dots$
- $A^0 \rightarrow \tau h^0 \dots$
- $h^0 \rightarrow \tilde{\chi}_1^0 \tilde{\chi}_1^0 \dots$

potentially large effects of radiative corrections

PRODUCTION CHANNELS:

Heinemeyer et al

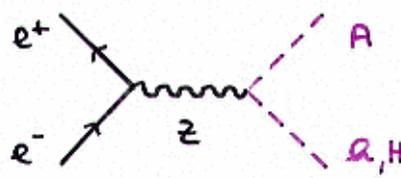
a) Higgs-strahlung:



$$\sigma_{Z h^0} = \sigma_0 \times \sin^2(\beta - \alpha)$$

$$\sigma_{Z H} = \sigma_0 \times \cos^2(\beta - \alpha)$$

b) assoc. production:



$$\sigma_{A h^0} = \sigma_0 \lambda \times \cos^2(\beta - \alpha)$$

$$\sigma_{A H} = \sigma_0 \lambda \times \sin^2(\beta - \alpha)$$

light h0 boson:

mutual \sin^2/\cos^2 in $Z h^0$ and $A h^0$ }
 \sin^2 small \rightarrow \cos^2 large, M_A small

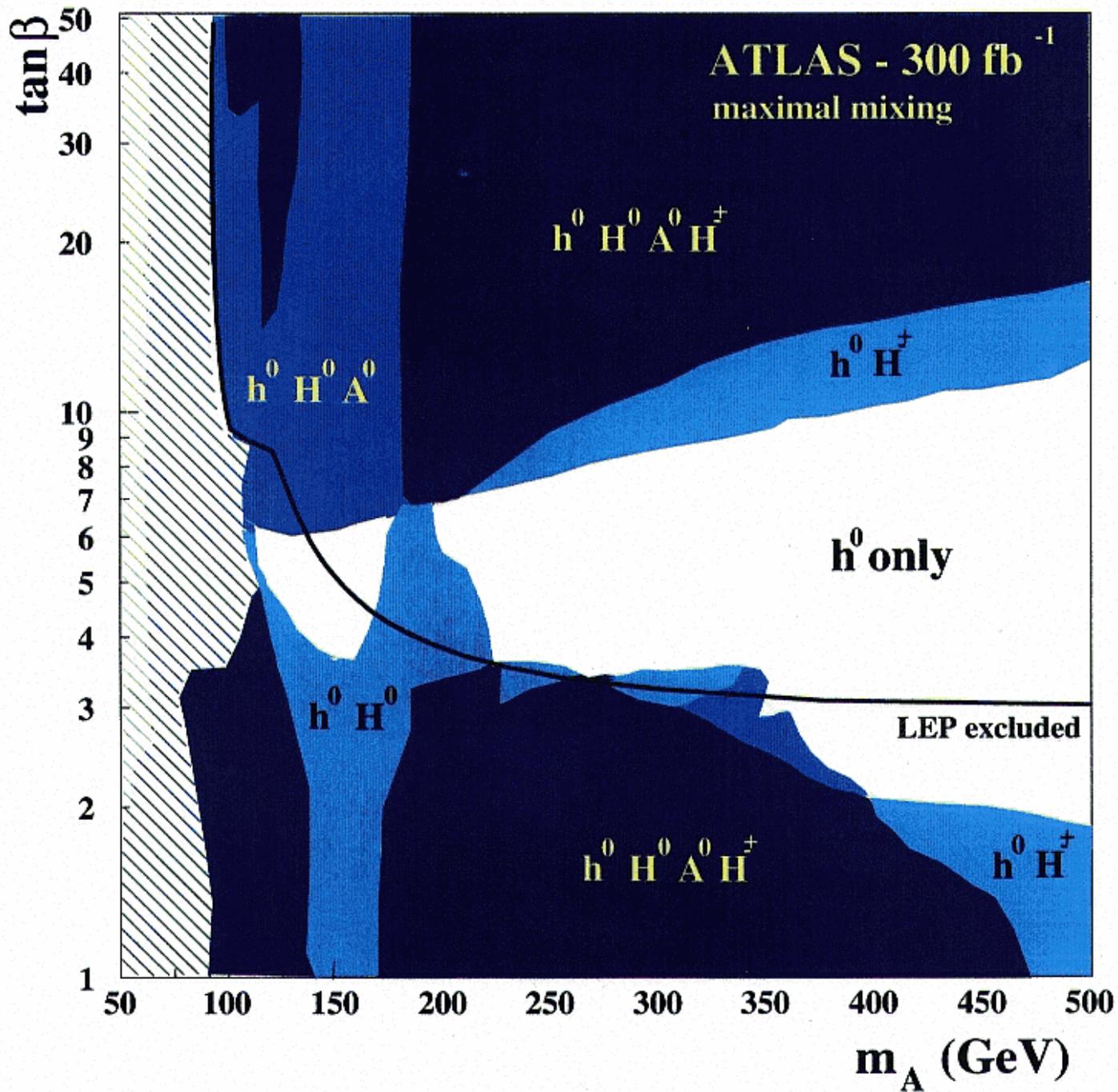
h^0 must be found / at ~ 250 GeV
 indep. of decay mode

heavy H^0, A^0, H^\pm :

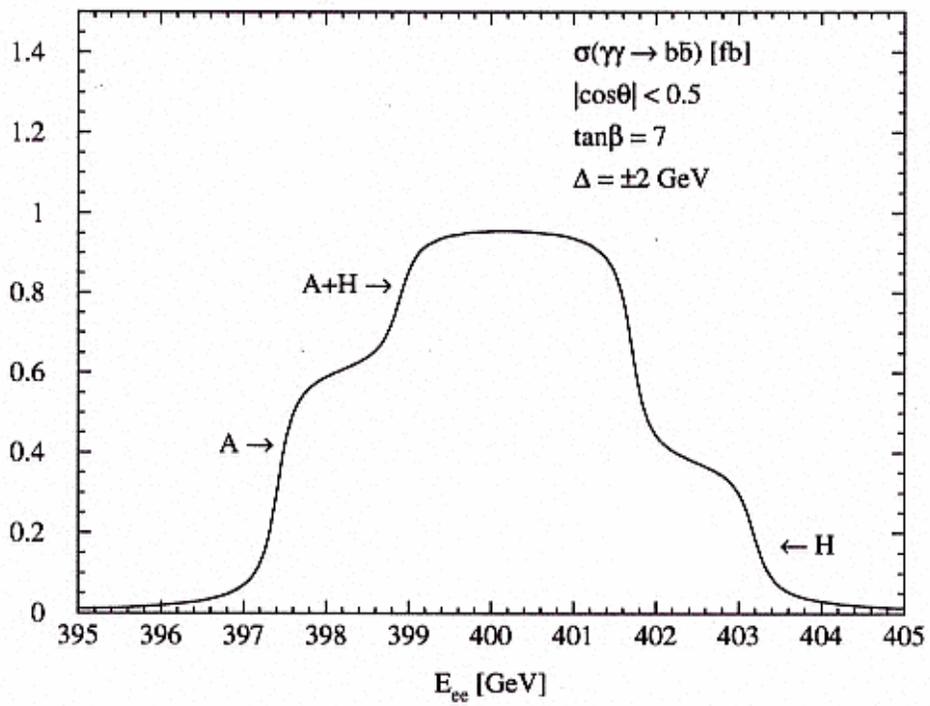
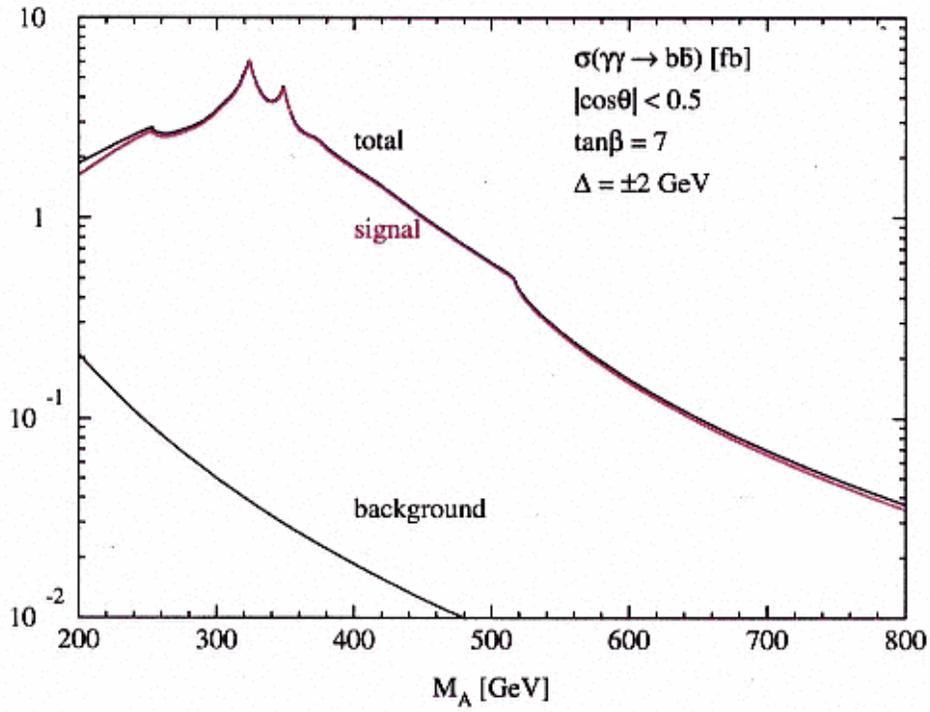
pair production in deeply regime

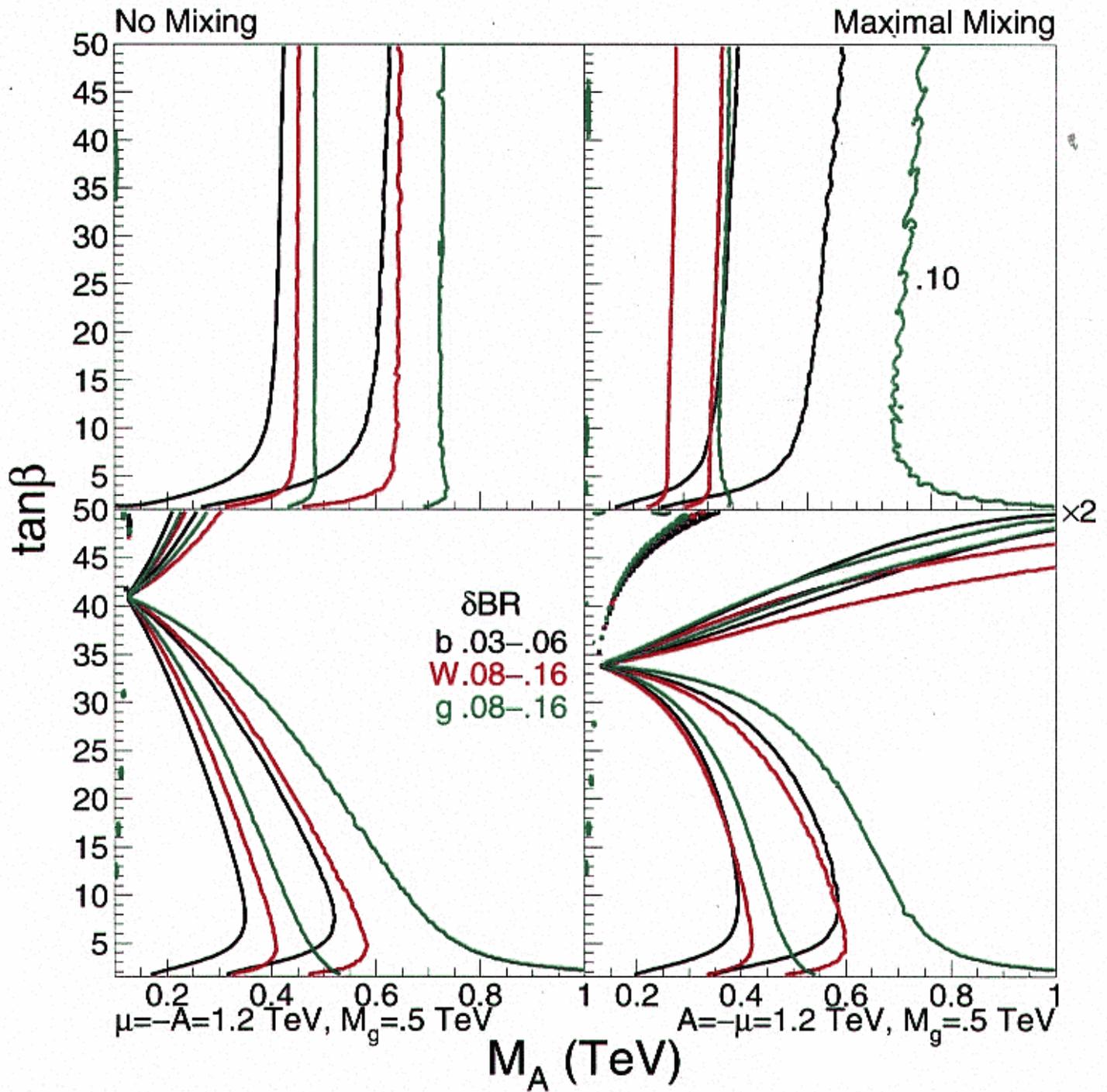
$$e^+e^- \rightarrow A^0 H^0, H^+ H^- \text{ with } M \lesssim \frac{1}{2} \sqrt{s}$$

CLIC: M up to 2.5 TeV



Müller
Kämer, Spica, Z





E. L. et al.

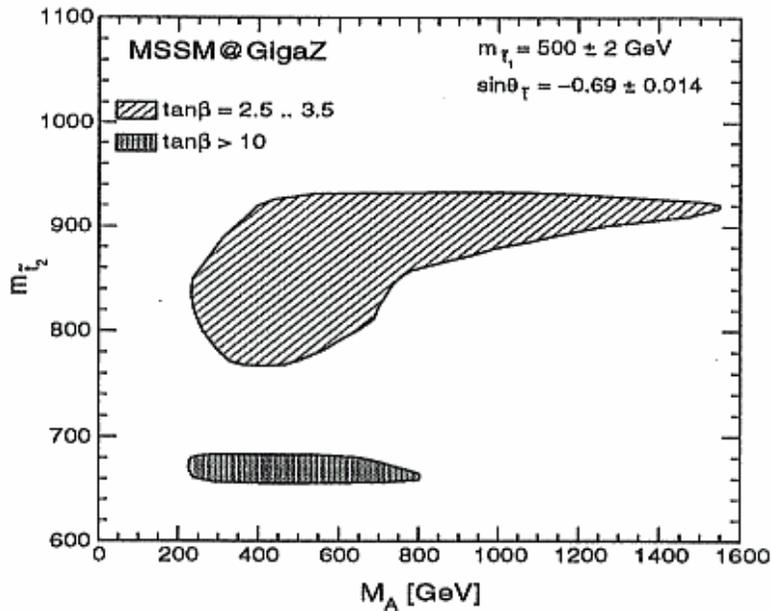
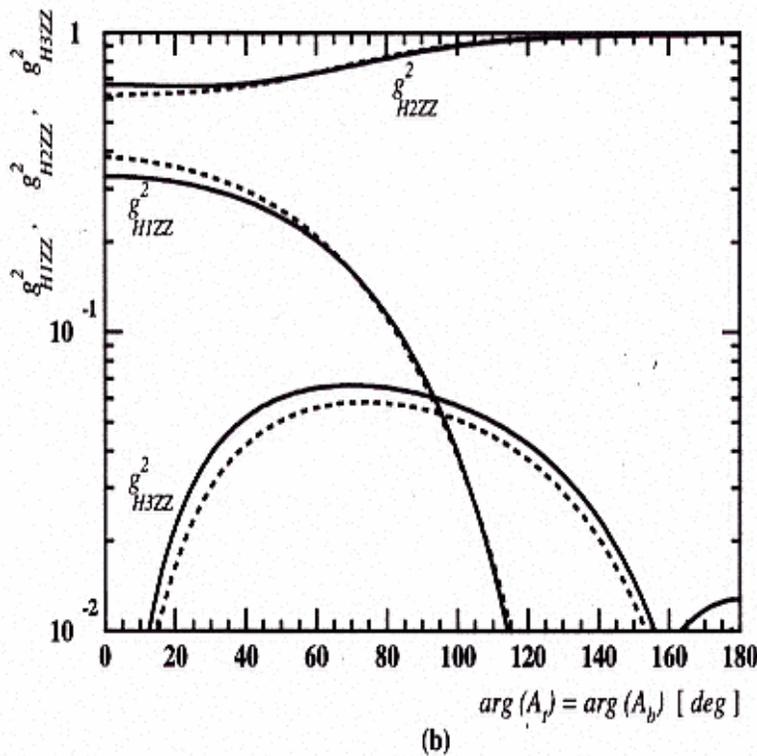
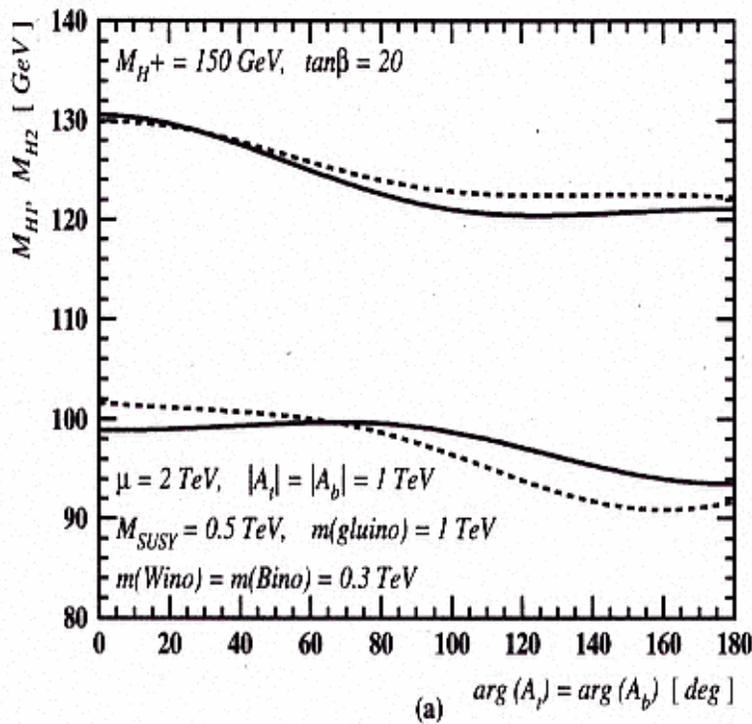


Figure 19: The region in the $m_A - m_{\tilde{t}_2}$ plane, allowed by 1σ errors obtained from the Giga-Z measurements of m_W and $\sin^2\theta_{\text{eff}}$: $m_W = 80.400 \pm 0.006 \text{ GeV}$, $\sin^2\theta_{\text{eff}} = 0.23140 \pm 0.00001$, and from the LC measurement of m_h : $m_h = 115 \pm 0.05 \text{ (exp.)} \pm 0.5 \text{ (theo.) GeV}$. $\tan\beta$ is assumed to be $\tan\beta = 3 \pm 0.5$ or $\tan\beta > 10$. The other parameters including their uncertainties are given by $m_{\tilde{t}_1} = 500 \pm 2 \text{ GeV}$, $\sin\theta_{\tilde{t}} = -0.69 \pm 0.014$, $A_b = A_t \pm 10\%$, $m_{\tilde{g}} = 500 \pm 10 \text{ GeV}$, $\mu = -200 \pm 1 \text{ GeV}$ and $M_2 = 400 \pm 2 \text{ GeV}$.

Couplings and Masses of Neutral Higgs ($\tan \beta = 20$)



(INDIRECT) ACCESS TO HEAVY MASSES :

note : LHC blind in wedge : centered at $\tan\beta \sim 8$
opening up at $M_A \sim 250 \text{ GeV}$ }

a) $\gamma\gamma \rightarrow H^0$ and A^0 : resonances in Compton Collides
 $M \leq 600 \text{ GeV}$ at LC of 800 GeV

b) $BR(h^0 \rightarrow b\bar{b})$ etc : $k_{bb} = -\sin\alpha / \cos\beta$: $\alpha \rightarrow M_A$

Okada
Battaglia

no mix : reach $\sim 700 \text{ GeV}$

Carone et

mixing : reach up to 1 to 2 TeV

[dependg on SUSY params]

c) Giga Z : $\sin^2\theta_{\text{eff}}, M_W$
 $m(h^0), M(\tilde{E}_i)$ and \tilde{E}_i } small $\tan\beta$: reach $\leq 1.6 \text{ TeV}$
large $\tan\beta$: $\leq 0.8 \text{ TeV}$

DISENTANGLING A^0 from H^0 :

decay : $H, A \rightarrow t\bar{t}$: $\frac{d\sigma}{d\varphi_{WW}} \sim 1 \mp \left(\frac{1}{4}\right)^2 \cos\varphi$

$\gamma\gamma$ prod : threshold scan

lin. pol. parallel : $\gamma\gamma_{\parallel} \rightarrow H^0$

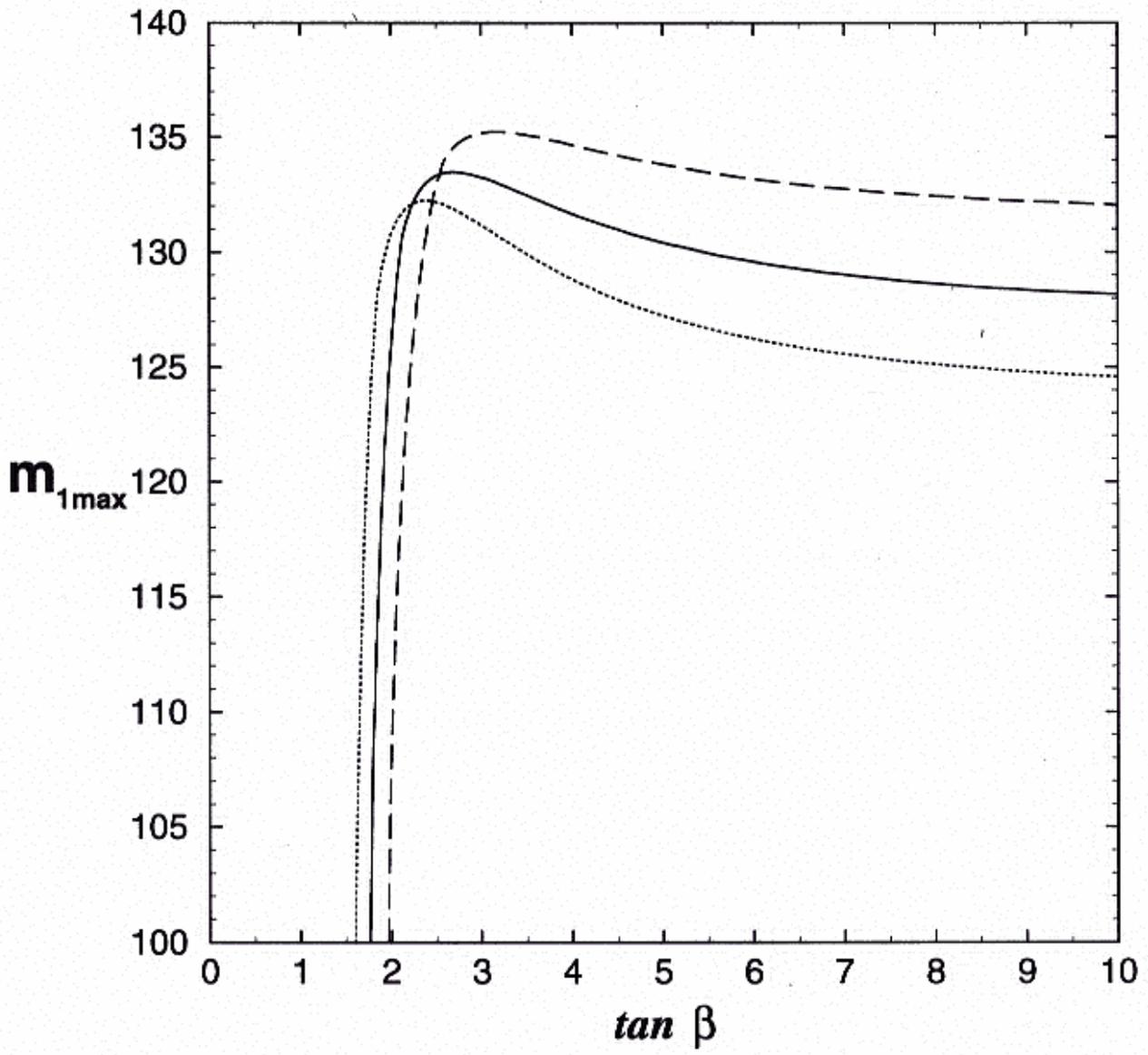
perpend. : $\gamma\gamma_{\perp} \rightarrow A^0$

CP : strongly affecting MSSM Higgs sector

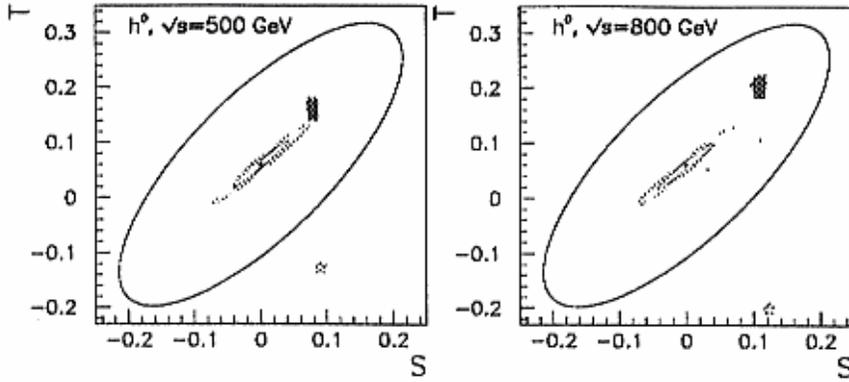
Carone et

$h^0, H^0, A^0 \rightarrow \text{CP mixtures } 1, 2, 3$

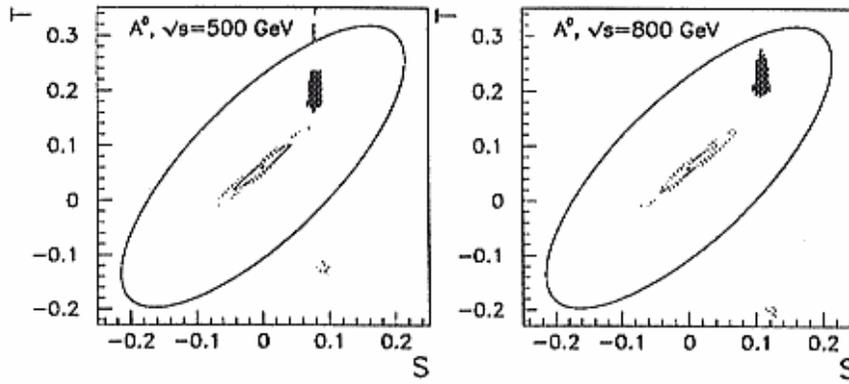
Ellwanger ea



S,T for $U=0$ and $\Delta\chi^2_{\min}$ in No-Discovery Zones



quion ea



(M+1)SSM : motivated by μ problem

Ellwanger

superstring theories ...

$$\begin{aligned} 2 \text{ doublets} + 1 \text{ triplet} : H^0, H^0 \oplus H^{10} &\Rightarrow H_1^0, H_2^0, H_3^0 \\ A^0 \oplus A^{10} &\Rightarrow A_1^0, A_2^0 \\ H^\pm &\Rightarrow H_1^\pm \end{aligned}$$

mass pattern \leftarrow Yukawa interaction $\lambda H_1 H_2 S$:

- λ bounded : $M_{H_1^0}^2 \leq M_Z^2 \cos^2 2\beta + \frac{2\lambda^2 H_Z^2}{g_1^2 + g_2^2} \sin^2 2\beta \rightarrow 135 \text{ GeV}$
- H^\pm below W^\pm : $M_{H^\pm}^2 = M_{W^\pm}^2 + M_A^2 - \lambda^2 v^2$

production :

- if $H_1^0 \sim S$ decpl $\Rightarrow H_2^0$ light \Rightarrow MSSM
- general no-lose theorem

Espeirsa
Junion

\uparrow not valid in non-SUSY 2HDM :

S, T, U analysis in GigaZ

Kauerzif
ea

REMARK :

Hall ea

Higgs pattern in general SUSY may be quite different from MSSM

Schub - Schwarz SUSY breaking \sim compactification of extra dim.

\rightarrow one light Higgs boson in dim = 4

$M(a) \leq 180 \text{ GeV} \oplus$ non-std SUSY properties

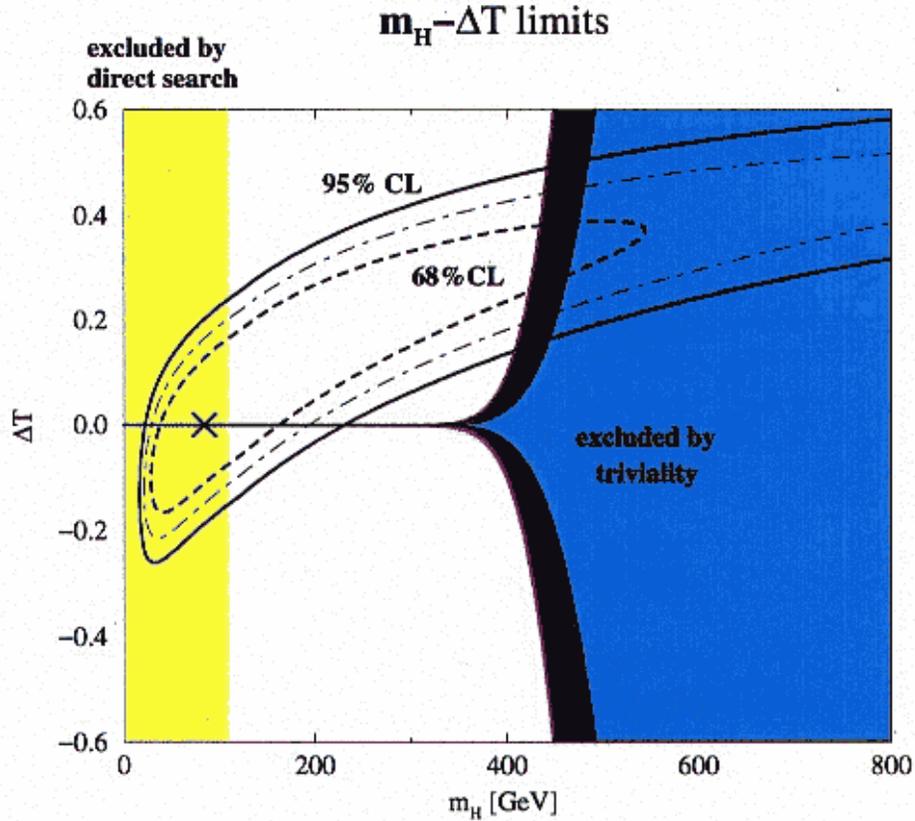


Figure 1: 68% and 95% CL bounds in the $(m_H, \Delta T)$ plane allowed by a fit to precision electroweak data [1, 2]. The best fit "standard model" value is shown by the cross on the $\Delta T = 0$ line. (Also shown by the dot-dash curve is the contour corresponding to $\Delta\chi^2 = 4$, whose intersection with the line $\Delta T = 0$ - at approximately 190 GeV - corresponds to the usual 95% CL upper bound quoted on the Higgs boson mass in the standard model.) The light region to the right is excluded by eqn. 1.3 for $b\kappa^2 = 4\pi$. The dark region denotes the additional area excluded for $b\kappa^2 = 4\pi^2$. The positive branches of the curves bounding these regions are lower bounds for ΔT in the top-seesaw and composite higgs models described in the text. Any $(m_H, \Delta T)$ with positive ΔT and to the left of the appropriate triviality curve can be realized in the corresponding model.

Chivukula et al

a) COMPOSITE HIGGS BOSON

Motivation: M_t large \Rightarrow t interrelated with eWSB?

top seesaw mechanism:

$$\oplus \underline{\chi \text{ singlet}}: \begin{bmatrix} t \\ b \end{bmatrix}_L \quad \begin{matrix} t_R \\ b_R \end{matrix} \quad \chi_L \quad \chi_R : \quad \varphi = \begin{bmatrix} \bar{\chi}_R & t_L \\ \bar{\chi}_R & b_L \end{bmatrix}$$

precision data Δg : $M_H \lesssim 500 \text{ GeV}$
 $M_\chi \gtrsim 5 \text{ TeV}$

$\oplus \chi, \omega$ singlets: spectrum of new scalar/pseudoscalar states,
 some with potentially small masses:

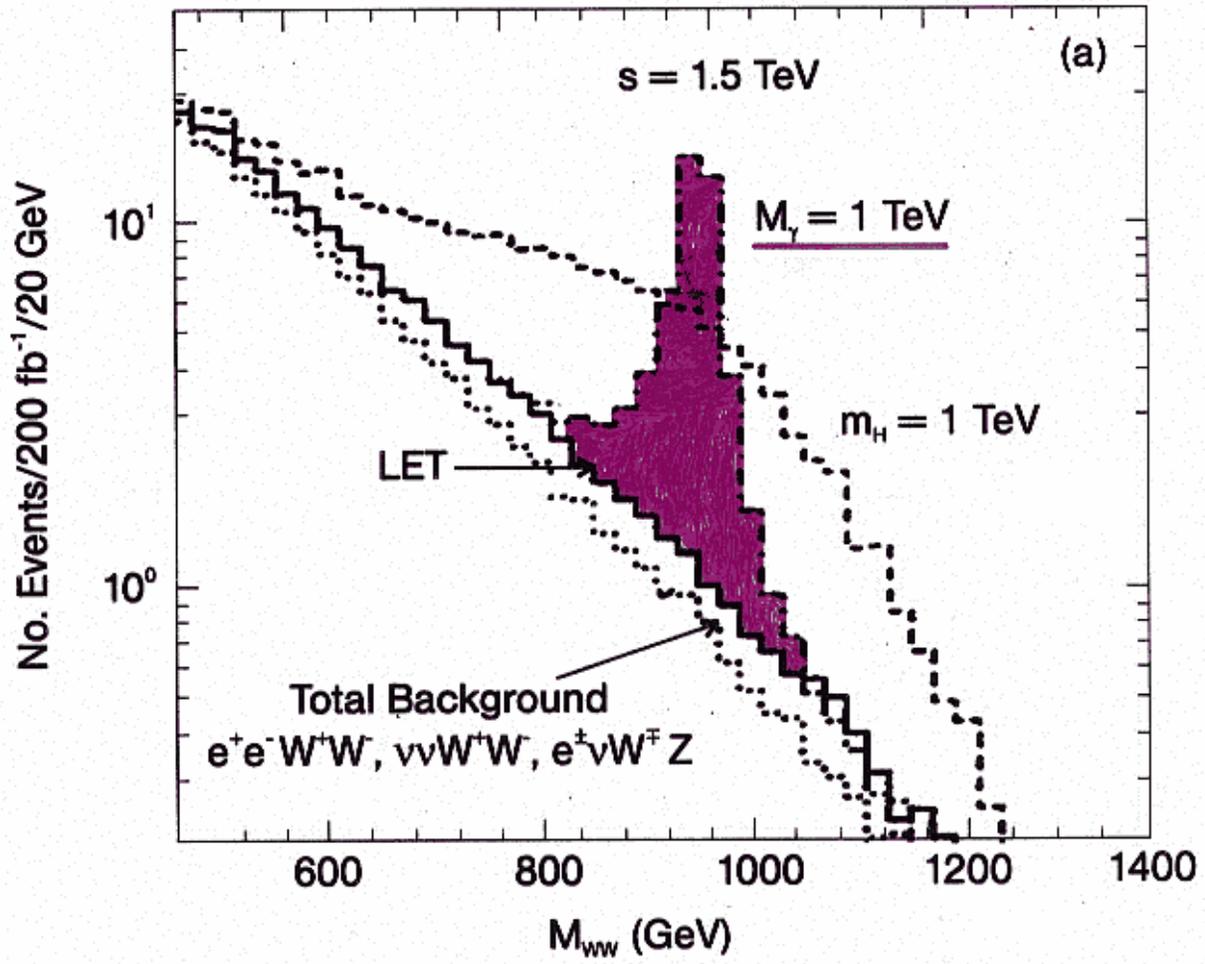
$$\left. \begin{array}{l} \text{light scalar} : h^0 \\ \text{light pseudoscalar} : A^0 \end{array} \right\}$$

$M_{h^0} > 2 M_{A^0}$: dominant h^0 decay: $h^0 \rightarrow A^0 A^0$

$$e^+e^- \rightarrow Z h^0 \rightarrow Z A^0 A^0 \rightarrow \gamma\gamma, 3\pi, \dots$$

Dobrescu et al

easy to detect at LC



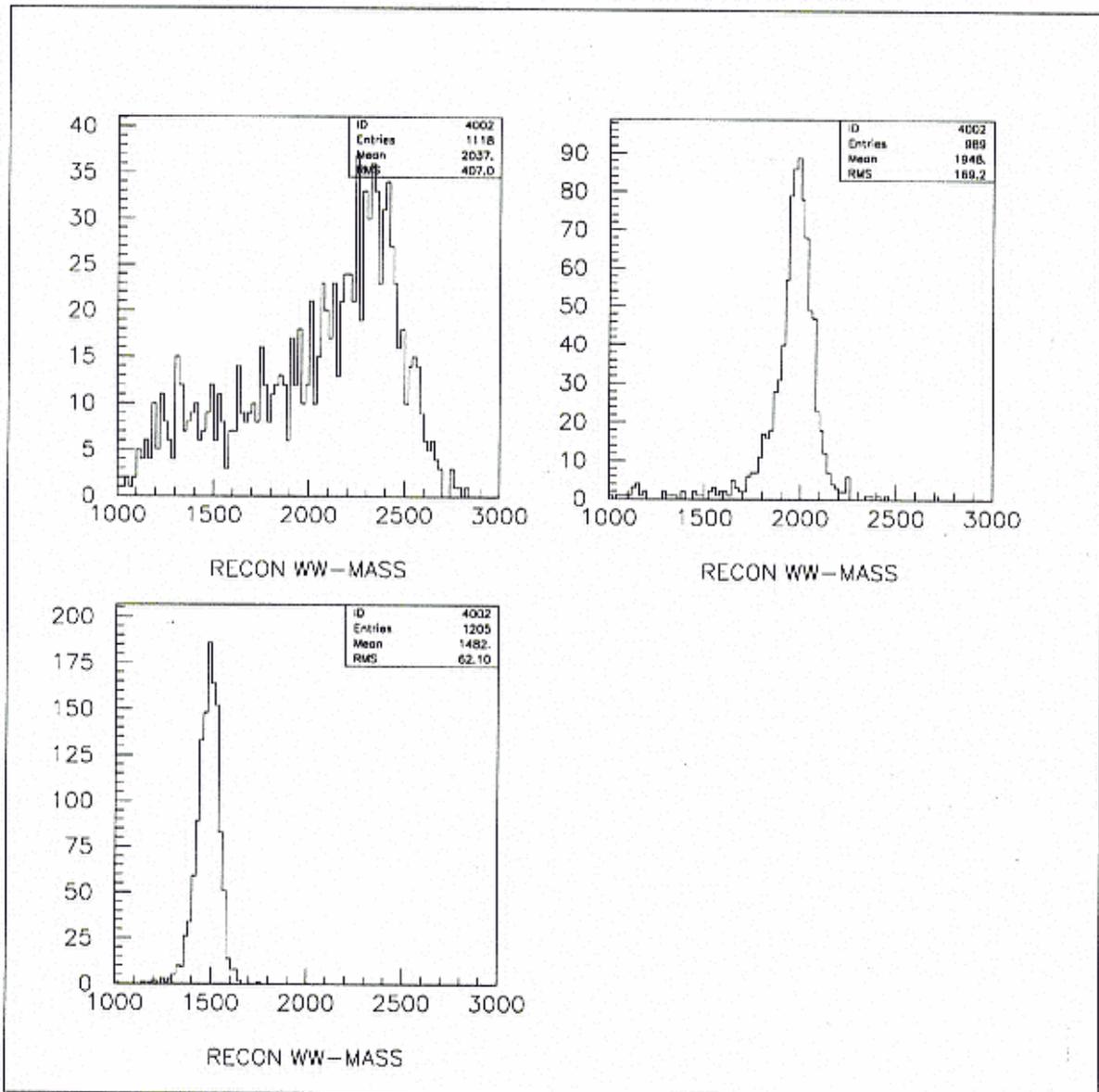
RECONSTRUCTED MASS

CLIC

$$M_{WW} = 1.5 \text{ TEV}, \Gamma = 35 \text{ GEV}$$

$$M_{WW} = 2 \text{ TEV}, \Gamma = 85 \text{ GEV}$$

$$M_{WW} = 2.5 \text{ TEV}, \Gamma = 250 \text{ GEV}$$



6) STRONGLY INTERACTING W BOSONS

no fundam. Higgs boson: strong WW interactions $\sim 1\text{TeV}$

← new strong interactions $\Lambda_* \lesssim 4\pi v \sim 3\text{TeV}$

global chiral symmetry: spontaneously broken
Goldstone bosons \sim longit. W states

:: deviations from SM predictions for energies below Λ_*
in precision analyses:

Barklow
Ollé et al
Casalbuoni

■ anom. bilin. qfg:

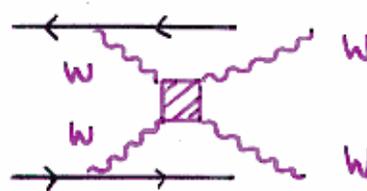


$\sqrt{s} = 500\text{ GeV}$

Λ_* up to 3 TeV

■ strong WW scattg:

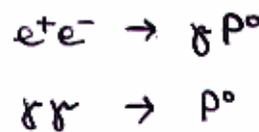
threshold below res.



$\sqrt{s} = 800\text{ GeV}$

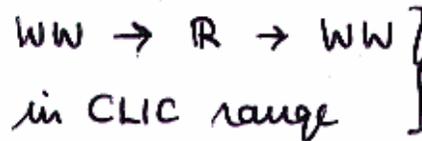
Λ_* up to 2.5 TeV

■ pseudo-Goldstone bosons:

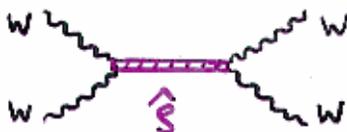


$M_0 \lesssim 400\text{ GeV}$
 $\lesssim 0.7\sqrt{s}$

:: resonance formation:



Barger et al
De Rooij



optimal environment for res. studies

5. CONCLUSIONS

- 1.) Essential elements of SM Higgs mechanism can be reconstructed in LC / phase I:
 - profile of Higgs boson
 - 3 lin. component of Higgs potential
- 2.) In supersymmetric theories, in LC / phase I access to
 - light Higgs boson and its properties
 - heavy Higgs bosons in e^+e^- mode : mass up to $\sqrt{s}/2$
in $\gamma\gamma$ mode : up to $\sim 0.7\sqrt{s_{ee}}$
and measurement of properties
 - no-lose theorem for SUSY Higgs discoveryin LC / phase II : access up to $\approx 2.5\text{ TeV}$
- 3.) Strong elw symmetry breaking and new strong interactions
 - covered in LC / phase I throughout threshold region
 - in LC / phase II : resonance formation covered

The central alternative scenarios, fundamental Higgs sector and strong elw symmetry breaking, can be reconstructed at e^+e^- linear colliders with unique completeness and precision.